

# Robust quantum inertial sensors with optical beam shaping

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## Introduction

### ABSTRACT/SUMMARY

Inertial sensors based on atom interferometry benefit from laser beams of homogeneous intensity and wavefront. Here we present a fiber-coupled athermalized beam shaper retaining these properties over an extended temperature range, developed for a field-deployed quantum gravimetry within the FIQUgS project. With optimally matched thermal properties of lens and housing materials, we obtain a top-hat collimator design with an RMS wavefront error of less than 6mλ over the entire temperature range and less than 3% intensity variation over the entire propagation range. Our results open perspectives for extending application of top-hat collimators towards different architectures of atomic inertial sensors demanding enhanced robustness, optical powers and laser beam propagation distances.

## Atom interferometry with top-hat beams

### BUILDING BLOCKS OF ATOM INTERFEROMETER ANALOGY WITH OPTICS

Optics		Matter waves
Coherent source	Plane wave/ laser field	De Broglie matter waves: $\lambda = h/p$ "monochromatic" source → laser cooling
Optics elements	Beam splitters, mirrors	"Atom optics" = Diffraction of atomic wave packets with lasers = Coherent interaction between laser field & an atom (e.g. Raman transition)

How does inertial sensitivity appear?  
 = Interferometric sequence (example of Mach-Zehnder)

Interferometric pattern:  
 $P_{|\vec{p}\rangle \rightarrow |\vec{p} + \hbar \vec{k}_{\text{eff}}\rangle} = \frac{1}{2}(1 - \cos \Delta\Phi)$

Classical equation of motion along the two paths  
 = Laser phase "imprinted" on diffracted wave function:  
 $\phi(t) = \vec{k}_{\text{eff}} \cdot \vec{r}(t)$   
 = Acceleration and rotation contributions  
 $\Delta\Phi = \vec{a} \cdot \vec{k}_{\text{eff}} T^2 - 2\vec{k}_{\text{eff}} \cdot (\vec{\Omega} \times \vec{v}) T^2$

### BENEFITS OF ATOM INTERFEROMETRY WITH TOP-HAT BEAMS

Non-homogeneous atom-light coupling across the laser beam:  
 $P(x) = \sin^2(\Omega_{\text{eff}}(x)\tau/2)$

Degraded efficiency of the atom-optics!

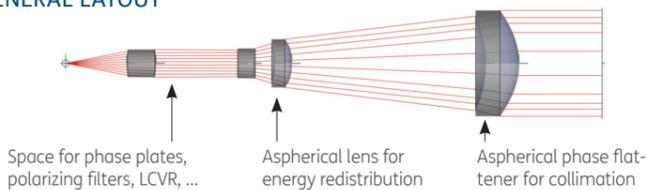
Fiber-coupled collimator with beam shaping optics  
 = Wavelength 780 nm  
 = Top-hat output profile 30 mm FWHM  
 = Beam uniformity 0.1  
 = RMS wavefront error <math>\lambda/10</math>  
 = Propagation length 2m  
 = Passive athermalization 0-50°C  
 = Housing titanium alloy

Gaussian beam  
 = Non-uniform transition probability + expansion of atomic cloud  
 → loss of interferometric contrast  
 = Spurious contributions to measured acceleration  
 → loss of measurement accuracy

Top-hat beam  
 = Uniform transition probability  
 = Elimination of spurious contributions

## Athermalization of fiber coupled beam shaper

### GENERAL LAYOUT



### CLASSICAL ATHERMALIZATION CONDITION

Collimation condition:  $\theta_1 = F(f_i(T_0), L_i(T_0)) = 0$  (1)  
 Athermalization condition:  $\frac{d\theta_1}{dT} \Big|_{T=T_0} = 0$  (2)  
 From (2):  $F(f_i(T_0), L_i(T_0)) \cdot \left( \alpha_H + \frac{dn/dT}{n-1} - \alpha_G \right) = 0$

Thermal power  $\gamma = \frac{dn/dT}{n-1} - \alpha_G$

Titanium:  $\alpha_H = 8.2 \cdot 10^{-6} \text{ K}^{-1}$   
 Fused silica:  $\gamma = 17.2 \cdot 10^{-6} \text{ K}^{-1}$   
 N-LAK21:  $\gamma = -8.3 \cdot 10^{-6} \text{ K}^{-1}$   
 N-KF9:  $\gamma = -10.1 \cdot 10^{-6} \text{ K}^{-1}$

Best matching glass: N-LAK21  
 Matching thermal properties of housing and glass

### ATHERMALIZATION OF BEAM SHAPING SYSTEMS

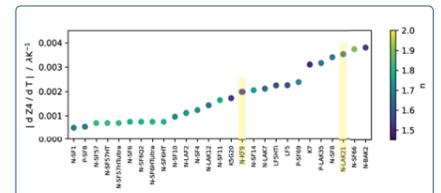
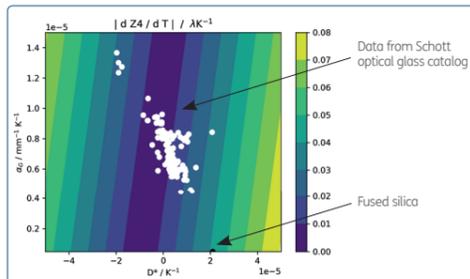
= Temperature change mainly introduces wavefront defocus (Z4)  
 = Optimization parameter  $\frac{dZ_4}{dT} = 0$   
 = Variation of glass properties

$$D^* = D_0 + \frac{E_0}{\lambda^2 - \lambda_k^2} \propto \frac{dn}{dT}$$

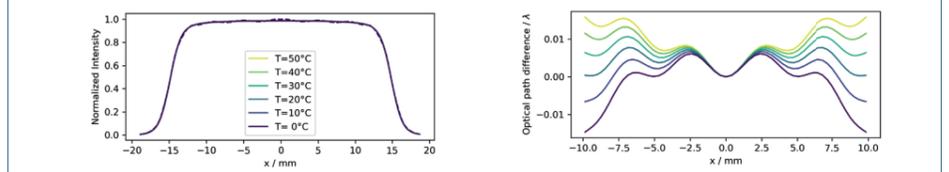
while keeping constant refractive index

Interpolation of  $|dZ_4/dT|$  to real glass data for selection of optimum glass with low refractive index and for titanium housing

Best matching glass with low refractive index: N-KF9  
 better athermalization than N-LAK21  
 $|dZ_4/dT|$  N-KF9 = 1.9mλ/K  
 $|dZ_4/dT|$  N-LAK21 = 3.5mλ/K



### Simulated performance of athermalized top-hat collimator with N-KF9



## Expected Gravimeter performance (FIQUgS project)

### THE DIFFERENTIAL QUANTUM GRAVIMETER IN DEVELOPMENT BY EXAIL

= Differential Quantum Gravimeter (DQG) is a dual interferometer quantum gravimeter  
 = Using a double tracking scheme, we are able to measure both gravity and vertical gravity gradient simultaneously  
 = Using the same laser for both interferometers ensures maximum common mode rejection on gradient  
 = Leverage the experience acquired in developing the AQG  
 = Gaussian beam collimator → top-hat collimator possible in future

Gravimeter  $\bar{g} = \frac{z_2 \theta_1 - z_1 \theta_2}{(z_2 - z_1)}$   
 Gradiometer  $\Gamma_{zz} = \frac{\theta_1 - \theta_2}{L}$

Perspective:  
 FIQUgS project aims deploying the dedicated version of DQG + spectral ground penetrating radar on the robotized platform for semi-autonomous geophysical surveys on the field

### SIMULATED TOP-HAT PROFILES AND INTERFEROMETRIC SEQUENCE

Simulation details:  
 = Atomic source rms expansion at 10 mm/ms  
 = T = 100 ms  
 = Top-hat profiles @ exact distances  
 = Consider no diffraction on apertures  
 = Counter-propagating pair of beams  
 = Raman coupling is given by:  
 $\Omega \propto E_{\text{forw}} * E_{\text{back}}^* = \sqrt{I_{\text{forw}}} * \sqrt{I_{\text{back}}} e^{i\Delta\phi}$

Coupling for the 1<sup>st</sup> laser pulse (t=0):  
 norm. amplitude  $|a|/|a|_{\text{mean}}$   
 relative phase  $\Delta\phi$  (λ)

Preliminary conclusion:  
 No significant contrast loss from residual inhomogeneity of top-hat laser beam profile